

Lecture 21

Grover search algorithm (extended with Qiskit examples)

of the course “Fundamentals of Quantum Computing”

(by  and **QUANTERALL**)

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PHYSICAL STUDIES



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Essence

Words as point in Hilbert space

Steps of the Grover search algorithm

Grover algorithm on a two-qubit system

Qiskit notebook

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Lov Kumar Grover

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The objects (haystack) must be represented as points in a Hilbert space. The vectors of the objects which we try to find span a subspace in this space.

The Grover algorithm constructs operators that transform a given initial state into a state which has a maximal component in the subspace of desired objects (*amplitude amplification*).

Definition 6.22 Let S denote the set of objects we are searching for, and let $m \geq 1$ be the cardinality of this set. The set S is called solution set, and we call its elements solutions. For the algorithm to search an $x \in S \subset \{0, \dots, N-1\}$, where $N = 2^n$, we define the input and output register as $\mathbb{H}^{I/O} = \mathbb{H}^{\otimes n}$. Furthermore, we denote the set of objects that are not a solution by

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$$S^\perp := \{0, \dots, N-1\} \setminus S$$

and define the subspaces

$$\begin{aligned}\mathbb{H}_S &:= \text{Span} \{ |x\rangle \mid x \in S \} \subset \mathbb{H}^{I/O} \\ \mathbb{H}_{S^\perp} &:= \text{Span} \{ |x\rangle \mid x \in S^\perp \} \subset \mathbb{H}^{I/O}\end{aligned}$$

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and the operators

$$\begin{aligned}P_S &:= \sum_{x \in S} |x\rangle \langle x| \\ P_{S^\perp} &:= \sum_{x \in S^\perp} |x\rangle \langle x| = \mathbf{1}^{\otimes n} - P_S\end{aligned}$$

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$$\begin{aligned}|\Psi_S\rangle &:= \frac{1}{\sqrt{m}} \sum_{x \in S} |x\rangle \\ |\Psi_{S^\perp}\rangle &:= \frac{1}{\sqrt{N-m}} \sum_{x \in S^\perp} |x\rangle\end{aligned}$$

Every state $|\Psi\rangle \in H^{I/O}$ can be decomposed as:

$$|\Psi\rangle = (P_{S^\perp} + P_S)|\Psi\rangle = \sum_{x \in S^\perp} \Psi_x |x\rangle + \sum_{x \in S} \Psi_x |x\rangle$$

Remember, measuring the state $|\Psi\rangle$ means:

Definition 5.35 Let $n \in \mathbb{N}$ and for $j \in \{0, \dots, n-1\}$ and $\alpha \in \{0, \dots, 3\}$ (or, equivalently, $\alpha \in \{0, x, y, z\}$) define

$$\Sigma_\alpha^j := \mathbf{1}^{\otimes n-1-j} \otimes \sigma_\alpha \otimes \mathbf{1}^{\otimes j} \in B_{sa}(\mathbb{H}^{\otimes n}),$$

where the σ_α are as in Definition 2.21. The **observation of a state in the quantum register** $\mathbb{H}^{\otimes n}$ is defined as the measurement of all compatible observables

$$\Sigma_z^j = \mathbf{1}^{\otimes n-1-j} \otimes \sigma_z \otimes \mathbf{1}^{\otimes j}$$

for $j \in \{0, \dots, n-1\}$ in the state of the quantum register. Such an observation is also called **read-out** or **measurement** of the register.

The goal of the algorithm to create states $|\Psi\rangle$ for which the probability to observe an $x \in S$ is maximized. This is accomplished by starting from an initial state $|\Psi_0\rangle$ and by applying suitable transformations which increase the component in H_S .

$$\begin{aligned} \mathbf{P} \left\{ \begin{array}{l} \text{Observation of } |\Psi\rangle \text{ projects} \\ \text{onto a state } |x\rangle \text{ with } x \in S \end{array} \right\} & \underbrace{=}_{(2.62)} \|P_S |\Psi\rangle\|^2 \underbrace{=}_{(6.137)} \left\| \sum_{x \in S} \Psi_x |x\rangle \right\|^2 \\ & \underbrace{=}_{(2.14)} \sum_{x \in S} |\Psi_x|^2 . \end{aligned}$$

$$g: \{0, \dots, N-1\} \longrightarrow \{0, 1\}$$

$$x \longmapsto g(x) := \begin{cases} 0 & \text{if } x \in S^\perp \\ 1 & \text{if } x \in S \end{cases}$$

$$\widehat{U}_g(|x\rangle \otimes |y\rangle) := |x\rangle \otimes |y \boxplus g(x)\rangle$$

where $|y\rangle$ belongs to an auxiliary register H^W .

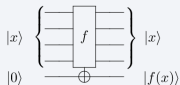
Reminder

$$|a \boxplus b\rangle := \bigotimes_{j=m-1}^0 |a_j \oplus^2 b_j\rangle$$

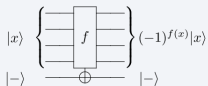
$$u \oplus^2 v := (u + v) \bmod 2$$

▼ Circuit Construction of a Grover Oracle (click to expand)

If we have our classical function $f(x)$, we can convert it to a reversible circuit of the form:



If we initialise the 'output' qubit in the state $|-\rangle$, the phase kickback effect turns this into a Grover oracle (similar to the workings of the Deutsch-Jozsa oracle):



We then ignore the auxiliary ($|-\rangle$) qubit.

Lemma 6.24 For the oracle U_g and the state

$$|\omega_i\rangle = |\omega_f\rangle = |-\rangle := \frac{|0\rangle - |1\rangle}{\sqrt{2}}$$

in the auxiliary register \mathbb{H}^W one has for arbitrary $|\Psi\rangle \in \mathbb{H}^{I/O}$

$$\hat{U}_g(|\Psi\rangle \otimes |-\rangle) = (R_{S^\perp}|\Psi\rangle) \otimes |-\rangle,$$

where

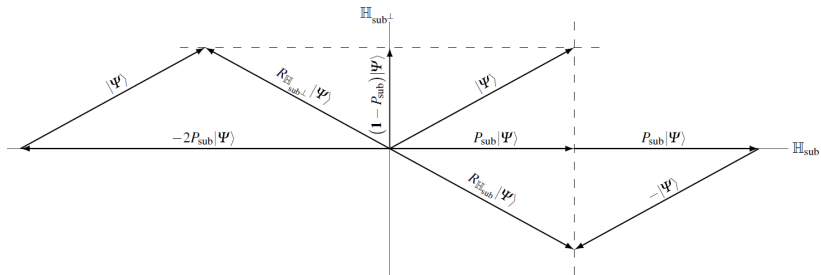
$$\begin{aligned} R_{S^\perp}|\Psi\rangle &= \sum_{x \in S^\perp} \Psi_x |x\rangle - \sum_{x \in S} \Psi_x |x\rangle \\ &= (\mathbf{1}^{\otimes n} - 2P_S)|\Psi\rangle. \end{aligned}$$

R_{S^\perp} can be viewed as reflection about H_{S^\perp} :

Definition 6.25 Let \mathbb{H}_{sub} be a subspace of the HILBERT space \mathbb{H} , and let P_{sub} be the projection onto this subspace. The **reflection about the subspace** \mathbb{H}_{sub} is defined as the operator

$$R_{\text{sub}} := 2P_{\text{sub}} - \mathbf{1}. \quad (6.147)$$

If the subspace is one-dimensional and spanned by a $|\Psi\rangle \in \mathbb{H}$, we simply write R_Ψ and call this a reflection about $|\Psi\rangle$.



Definition 6.26 Let S be the solution set with cardinality $m \geq 1$. For the algorithm to search for an $x \in S \subset \{0, \dots, N-1\}$, where $N = 2^n$, we define the initial state in the input/output register as

$$|\Psi_0\rangle := \frac{1}{\sqrt{N}} \sum_{x=0}^{N-1} |x\rangle \in \mathbb{H}^{I/O}. \quad (6.148)$$

Moreover, we define the angle

$$\theta_0 := \arcsin \left(\sqrt{\frac{m}{N}} \right) \in \left[0, \frac{\pi}{2} \right], \quad (6.149)$$

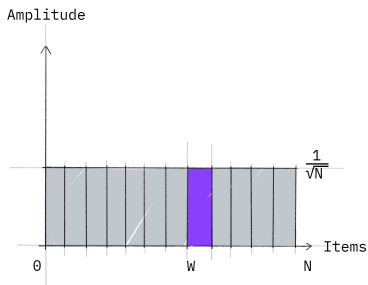
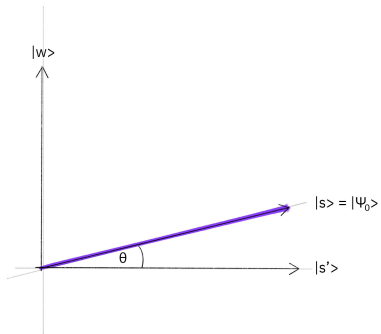
and with the help of $|\Psi_0\rangle$ we define the operator

$$R_{\Psi_0} = 2|\Psi_0\rangle\langle\Psi_0| - \mathbf{1}^{\otimes n} \quad (6.150)$$

on $\mathbb{H}^{I/O}$ as well as the initial state in the composite system

$$|\widehat{\Psi}_0\rangle := |\Psi_0\rangle \otimes |-\rangle \in \mathbb{H}^{I/O} \otimes \mathbb{H}^W. \quad (6.151)$$

$$|\Psi_0\rangle = \cos \theta_0 |\Psi_{S^\perp}\rangle + \sin \theta_0 |\Psi_S\rangle$$



$$|s'\rangle \equiv |\Psi_{S\perp}\rangle, |w\rangle \equiv |\Psi_S\rangle$$

Definition 6.27 The **GROVER iteration** is defined as the operator

$$\widehat{G} := (R_{\Psi_0} \otimes \mathbf{1}) \widehat{U}_g$$

on $\mathbb{H}^{I/O} \otimes \mathbb{H}^W$.

As we will now show, the GROVER iteration G transforms separable states in $\mathbb{H}^{I/O} \otimes \mathbb{H}^W$ of the form $|\widehat{\Psi}_j\rangle = |\Psi_j\rangle \otimes |-\rangle$ to separable states $|\widehat{\Psi}_{j+1}\rangle = |\Psi_{j+1}\rangle \otimes |-\rangle$ of a similar form. We will see that in the input/output register $\mathbb{H}^{I/O}$ an application of \widehat{G} can then be viewed as a rotation of $2\theta_0$ in $\mathbb{H}^{I/O}$ in the direction of $|\Psi_S\rangle$.

Proposition 6.28 For $j \in \mathbb{N}_0$ let

$$|\widehat{\Psi}_j\rangle := \widehat{G}^j |\widehat{\Psi}_0\rangle.$$

Then we have for all $j \in \mathbb{N}_0$

$$|\widehat{\Psi}_j\rangle = |\Psi_j\rangle \otimes |-\rangle$$

with $|\Psi_j\rangle \in \mathbb{H}^{I/O}$ and

$$|\Psi_j\rangle = \cos \theta_j |\Psi_{S^\perp}\rangle + \sin \theta_j |\Psi_S\rangle,$$

where

$$\theta_j = (2j+1)\theta_0.$$

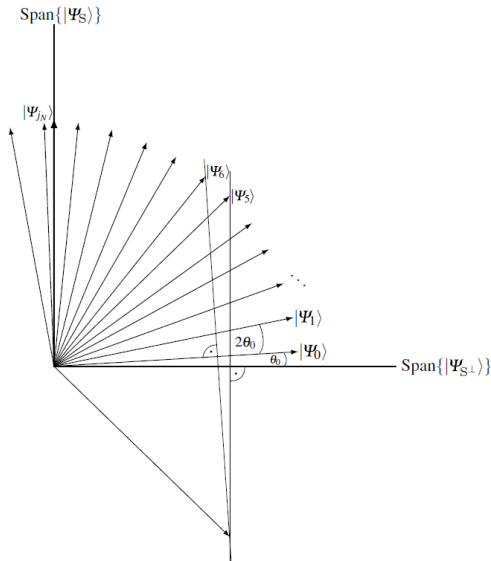
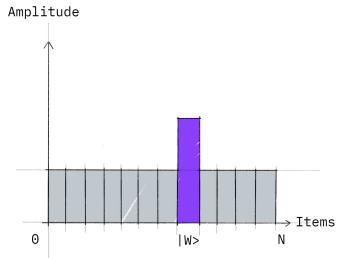
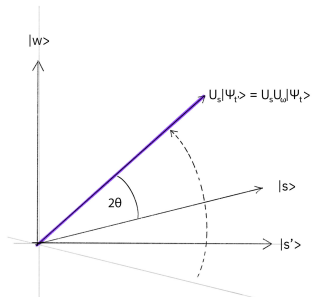
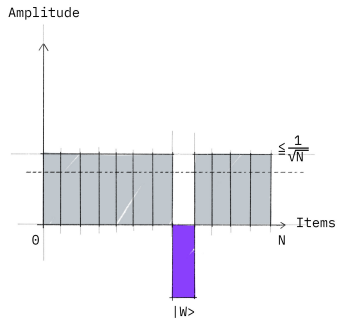
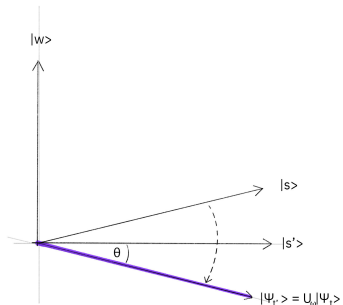


Fig. 6.6 Geometry of the GROVER iteration in the input/output register with $m = 5, N = 2^{10}$ and $j_N = 11$. In the two-dimensional subspace $\text{Span}\{|\Psi_{S\perp}\rangle, |\Psi_S\rangle\}$ the initial state $|\Psi_0\rangle$ is rotated towards $|\Psi_S\rangle$. The illustrated transition from $|\Psi_5\rangle$ to $|\Psi_6\rangle$ shows that \hat{G} in the sub-system I/O first performs a reflection about $|\Psi_{S\perp}\rangle$ and then a reflection about $|\Psi_0\rangle$. The vector immediately to the right of $|\Psi_{j_N}\rangle$ is $|\Psi_S\rangle$ and is a state in the subspace \mathbb{H}_S of the solution set. We can see that $|\Psi_{j_N}\rangle$ comes close to that



$$\mathbf{P} \left\{ \begin{array}{l} \text{Observation of } |\Psi_j\rangle \text{ projects} \\ \text{onto a state } |x\rangle \text{ with } x \in S \end{array} \right\} = \|P_S |\Psi_j\rangle\|^2 \underbrace{=}_{\substack{(6.137), (6.138), \\ (6.154)}} \sin^2 \theta_j$$

Lemma 6.29 *Let S be the solution set with cardinality $m \geq 1$, and let $N = 2^n$ be the number of objects in which we search for solutions. If we apply the GROVER iteration \hat{G}*

$$j_N := \left\lfloor \frac{\pi}{4 \arcsin(\sqrt{\frac{m}{N}})} \right\rfloor \quad (6.158)$$

times to $|\hat{\Psi}_0\rangle$ and observe the state $|\Psi_{j_N}\rangle$ in the input/output register, then the probability to observe in the sub-system $\mathbb{H}^{I/O}$ a state $|x\rangle$ with $x \in S$ satisfies

$$\mathbf{P} \left\{ \begin{array}{l} \text{Observation of } |\Psi_j\rangle \text{ projects} \\ \text{onto a state } |x\rangle \text{ with } x \in S \end{array} \right\} \geq 1 - \frac{m}{N}. \quad (6.159)$$

Steps of the Grover search algorithm

Input: A set $\{0, \dots, N-1\}$ of $N = 2^n$ objects that contains a subset S of $m \geq 1$ objects to be searched for and an oracle-function $g : \{0, \dots, N-1\} \rightarrow \{0, 1\}$ that takes the value 1 in S and the value 0 elsewhere

Step 1: In $\mathbb{H}^{I/O} \otimes \mathbb{H}^W = \mathbb{H}^{\otimes n} \otimes \mathbb{H}$ prepare the composite system in the state $|\hat{\Psi}_0\rangle = |\Psi_0\rangle \otimes |-\rangle$ with

$$|\Psi_0\rangle = \frac{1}{\sqrt{N}} \sum_{x=0}^{N-1} |x\rangle.$$

The number of computational steps required for Step 1 scales for $N \rightarrow \infty$ with

$$S_{\text{GROVER1}}(N) \in O(1)$$

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Step 2: With $\theta_0 = \arcsin(\sqrt{\frac{m}{N}})$ apply the transform $\hat{G} = (R_{\Psi_0} \otimes \mathbf{1})\hat{U}_g$

$$j_N = \left\lfloor \frac{\pi}{4\theta_0} \right\rfloor$$

times to $|\hat{\Psi}_0\rangle$ in order to transform the composite system to the state

$$|\hat{\Psi}_{j_N}\rangle = \hat{G}^{j_N} |\hat{\Psi}_0\rangle.$$

The number of computational steps required for Step 2 scales for $N \rightarrow \infty$ with

$$S_{\text{GROVER2}}(N) \in O\left(\sqrt{\frac{N}{m}}\right)$$

Step 3: Observe the sub-system $\mathbb{H}^{I/O}$ and infer from the observed state $|x\rangle$ the value $x \in \{0, \dots, N-1\}$. The number of computational steps required for Step 3 scales for $N \rightarrow \infty$ with

$$S_{\text{GROVER3}}(N) \in O(1)$$

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Step 4: By evaluating $g(x)$, check if $x \in S$. The number of computational steps required for Step 4 scales for $N \rightarrow \infty$ with

Output: A solution $x \in S$ with probability no less than $1 - \frac{m}{N}$

Grover algorithm on a two-qubit system

in this case Eq. 6.149 becomes

$$\theta_0 = \arcsin(\sqrt{1/4}) = \pi/6$$

j_N for the equation 6.158 is

$$j_N = \left\lceil \frac{\pi}{4 \arcsin(\sqrt{1/4})} \right\rceil = 1$$

Then the Grover transformation $G^{j_N} \equiv G^1$ yields

$$\theta_{j_N} = (2j_N + 1)\theta_0 = \pi/2, \text{ i.e.}$$

$$G^{j_N}|\Psi_0\rangle = \cos\theta_{j_N}|\Psi_{S_\perp}\rangle + \sin\theta_{j_N}|\Psi_S\rangle \equiv |\Psi_S\rangle$$

For the case when $S \equiv |11\rangle$ see the notebook linked in the following slide.

Qiskit notebook

<https://qiskit.org/textbook/ch-algorithms/grover.html>

N E X T L E C T U R E

N O V E M B E R 25, 2022

THANK YOU FOR
YOUR ATTENTION!

БЛАГОДАРЯ ЗА
ВНИМАНИЕТО!